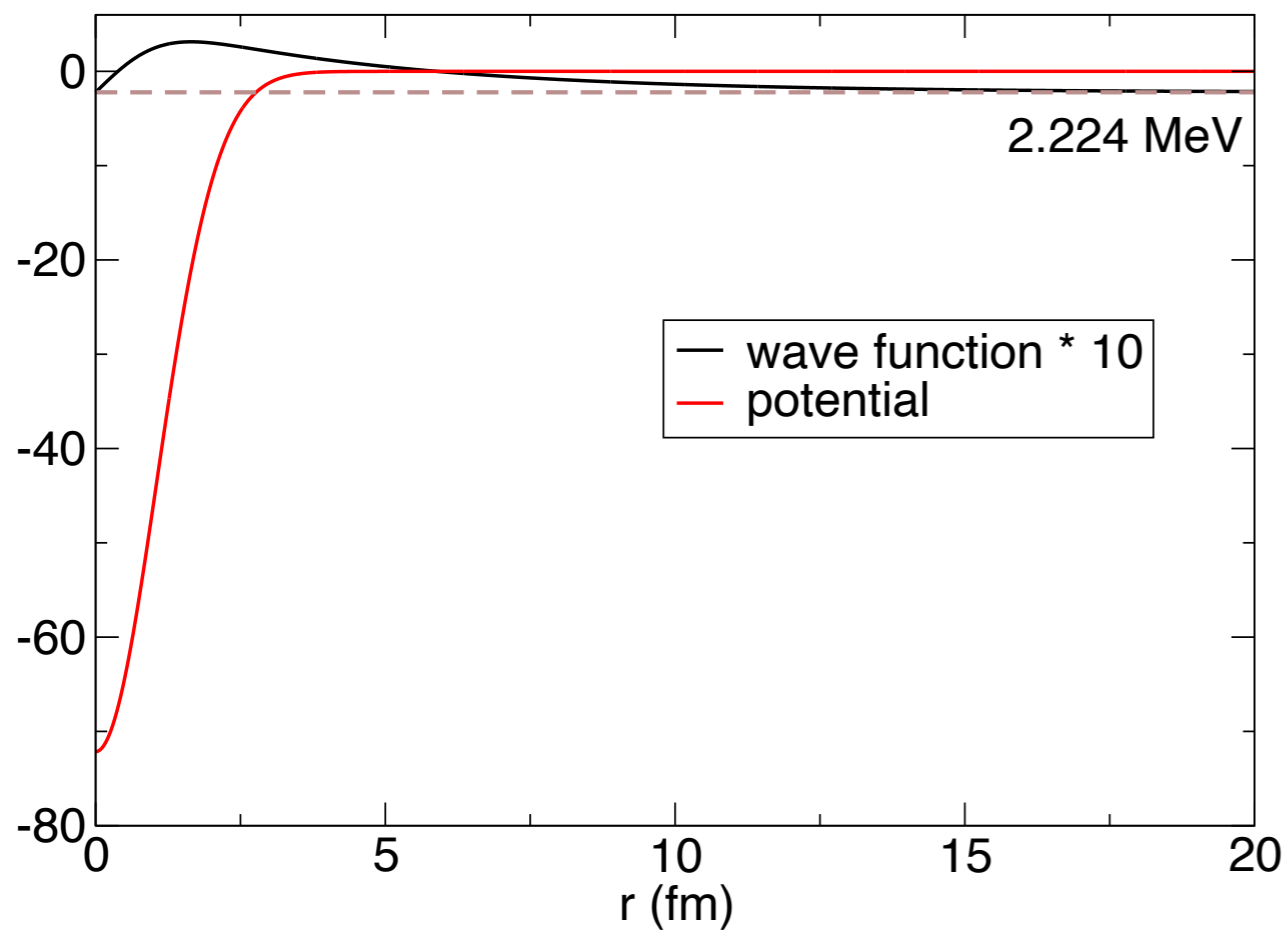
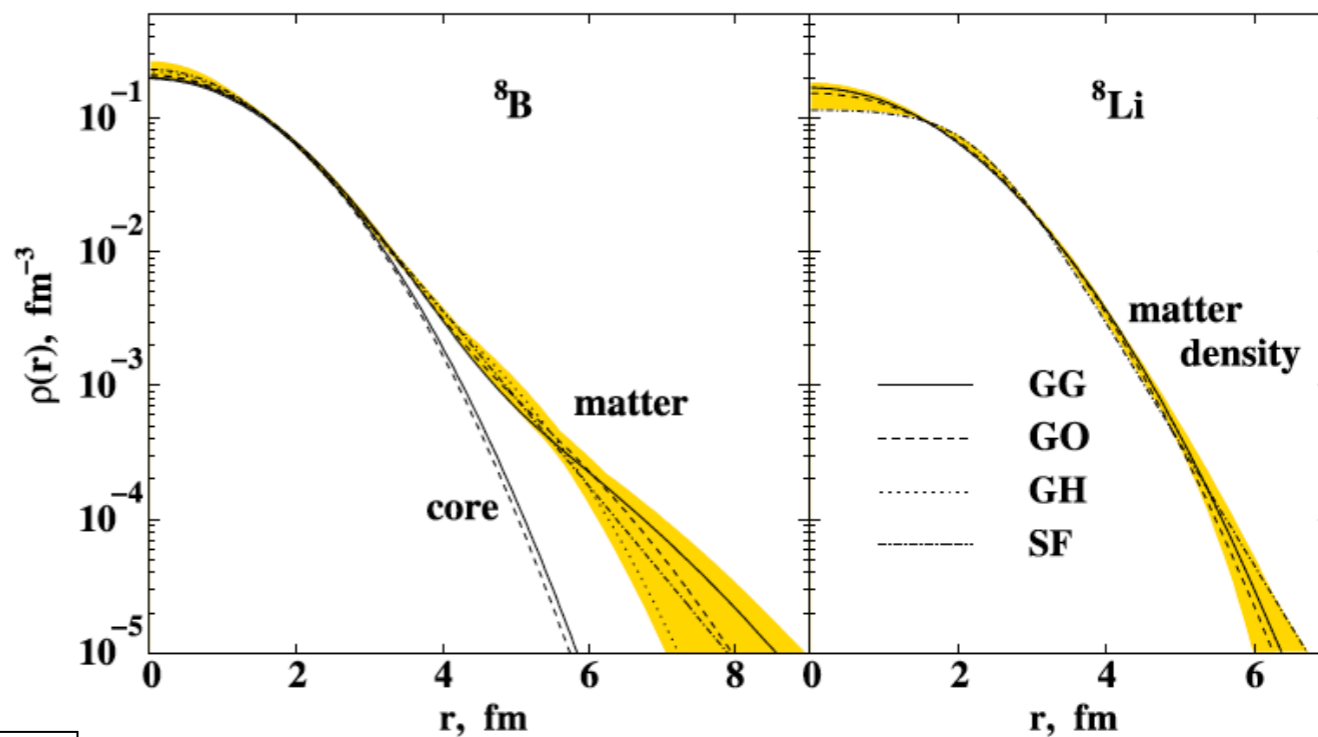




在动量表象下求解 $n-p$ 束缚态

量子多体角度
密度分布



量子少体角度
波函数与势相对空间分布

角动量基下的傅里叶变换

束缚态波函数是束缚态在坐标表象下的投影

上周求解的是径向波函数

$$\phi(\vec{r}) = \langle \vec{r} | \phi \rangle$$

$$\phi_l^m(r) = \langle rlm | \phi \rangle$$

同样的束缚态可以投影到动量表象下

$$\phi(\vec{k}) = \langle \vec{k} | \phi \rangle$$

通过傅里叶变化可以实现坐标表象与动量表象之间的变换

归一化条件

$$\langle \vec{r} | \vec{k} \rangle = \frac{1}{N_1} (2\pi)^{-3/2} e^{i\vec{k}\vec{r}}$$

$$\int d^3k e^{i\vec{k}(\vec{r}-\vec{r}')} = (2\pi)^3 \delta(\vec{r}-\vec{r}')$$

$$\langle \vec{k}' | \vec{k} \rangle = N_2 \delta(\vec{k}' - \vec{k})$$

$$\langle \vec{r}' | \vec{r} \rangle = N_3 \delta(\vec{r}' - \vec{r})$$

N_1, N_2, N_3 满足

$$N_1^2 N_2 N_3 = 1$$

我们选取

$$N_1 = N_2 = N_3 = 1$$

角动量基下的傅里叶变换

分波展开

$$\langle \vec{r} | \vec{k} \rangle = (2\pi)^{-3/2} e^{i\vec{k}\vec{r}} = \sum_{lm} \sqrt{\frac{2}{\pi}} \frac{1}{kr} i^l F_l(kr) Y_{lm}^*(\hat{k}) Y_{lm}(\hat{r})$$

↓
Coulomb wave function

Spherical Bessel函数与Coulomb wave函数的关系?

角动量基下的傅里叶变换

$$\langle rlm | klm \rangle = \sqrt{\frac{2}{\pi}} \frac{1}{kr} i^l F_l(kr)$$

动量表象下的势

坐标表象下的势

$$V(r) = V_0 \exp(-r^2/a^2) \quad a = 1.484$$

上周我们解得

$$V_0 = -72.167 \text{ MeV}$$

实际上“势”是算符

$$\langle r'lm | V | rlm \rangle = \frac{\delta(r - r')}{r^2} V(r)$$

↖ ↗ ↖
非定域 定域

动量表象下的势

$$\begin{aligned} \langle k'lm | V | klm \rangle &= \int_0^\infty r^2 r'^2 dr dr' \langle k'lm | r'lm \rangle \langle r'lm | V | rlm \rangle \langle rlm | klm \rangle \\ &= \int_0^\infty r^2 r'^2 dr dr' \langle k'lm | r'lm \rangle \frac{\delta(r - r')}{r^2} V(r) \langle rlm | klm \rangle \\ &= \frac{2}{\pi} \frac{1}{k'k} \int_0^\infty F_l(k'r) V(r) F_l(kr) dr \end{aligned}$$

动量空间下求解np束缚态

束缚态薛定谔方程

$$(E - H)|\phi\rangle = 0$$

变换得

$$(E - T)|\phi\rangle = V|\phi\rangle \quad \longrightarrow \quad |\phi\rangle = \frac{1}{E - T}V|\phi\rangle$$

投影到动量空间下

$$\langle klm | \phi \rangle = \int_0^\infty \frac{1}{E - \frac{(\hbar k)^2}{2\mu}} V_l(k, k') \langle k'lm | \phi \rangle k'^2 dk'$$

取 $\hbar k = k$ 注意单位的变换

积分运算在数值运算中为求和运算

$$\phi(k_i) = \sum_j \left(k_j^2 \omega_j \frac{1}{E - \frac{k_i^2}{2\mu}} V_l(k_i, k_j) \right) \phi(k_j)$$

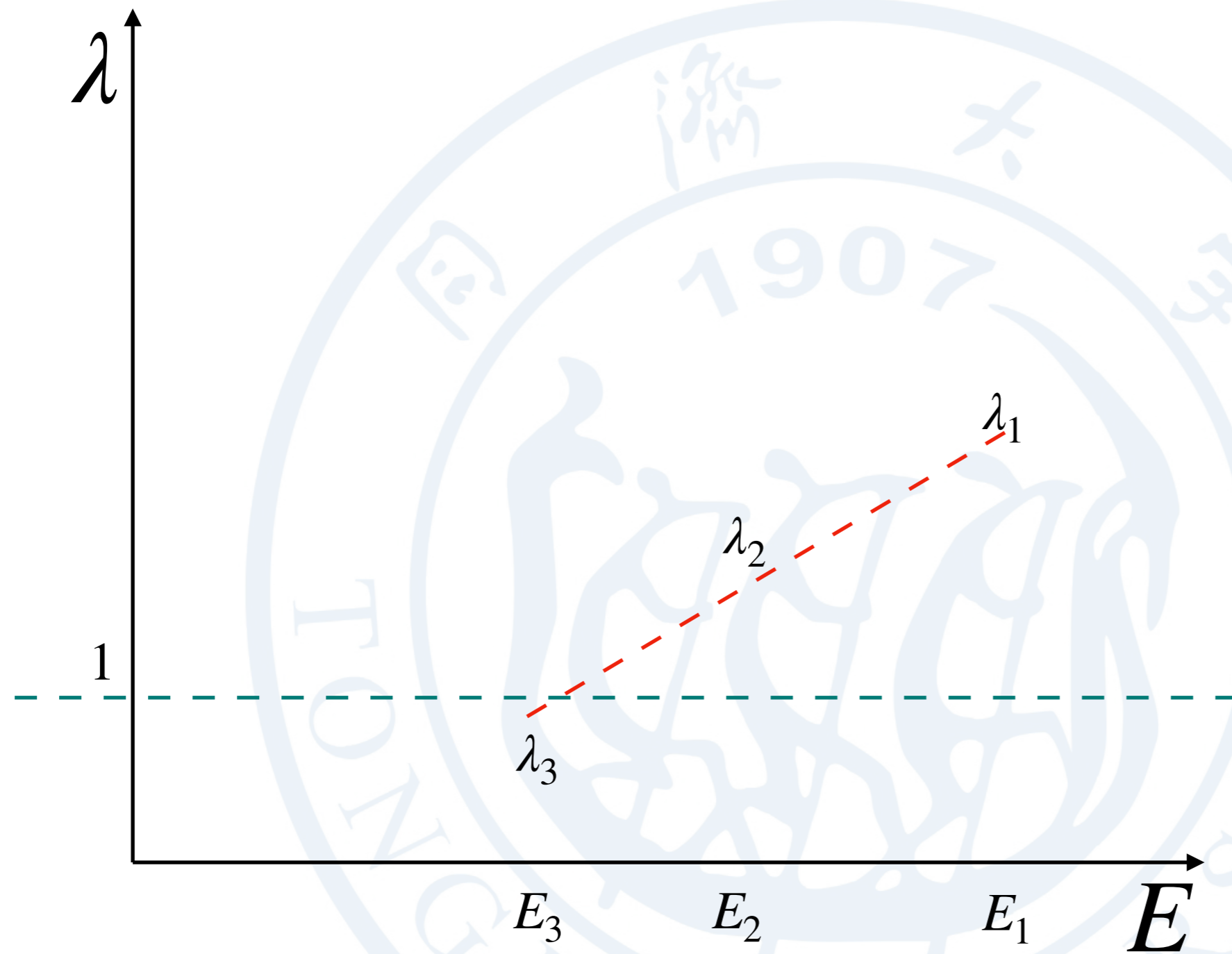
A_{ij}

求解束缚态变成求解本征值问题

$$\lambda \phi = \mathbf{A} \phi$$

当E值合适时 $\lambda = 1$

Secant method



迭代下去直到

$|E_n - E_{n-1}|$ 足够小

$$E_3 = E_2 - \frac{E_1 - E_2}{\lambda_1 - \lambda_2} (\lambda_2 - 1)$$

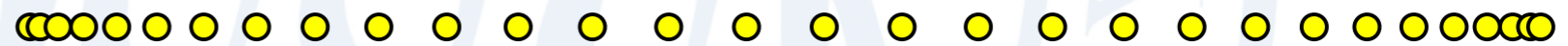
积分变求和

$$\int f(x)dx \approx \sum_i f(x_i)\omega_i$$

Simpson Map



Gaussian Map



np1

np2

TRNS Map



计算用单位

为了方便起见，把MeV统一变换成 fm^{-1}

通过 $hc = 197.3269718 \text{ MeV} \cdot fm$

比如中子的质量为 939.5983 MeV ，那么换算后为 4.7616 fm^{-1}

