Basics of Direct Nuclear Reaction<br>junzliu<br>April 2022

## 1 Big picture of a direct reaction

### 1.1 Reaction

The reaction takes the form:

$$
\begin{equation*}
a+A \rightarrow c+C \tag{1}
\end{equation*}
$$

where a is called the projectile and A is called the target. The character c stands for 'channel'. Channel is a key concept in direct nuclear reaction. The interaction between two nuclei can produce different partitions of the a+A system. Besides, the internal state may be excited. These are all different channels of the reaction.

### 1.2 Notations

The total Hamiltonian of the system is written as:

$$
\begin{equation*}
H=H_{A}+H_{a}+T_{\alpha}+V_{\alpha} \tag{2}
\end{equation*}
$$

where $\alpha$ stands for the incident channel.
The internal states are represented by:

$$
\begin{equation*}
H_{A} \Phi_{A}=\epsilon_{A} \Phi_{A}, \quad H_{a} \Phi_{a}=\epsilon_{a} \Phi_{a} \tag{3}
\end{equation*}
$$

But the Hamiltonian of the whole system can be expressed in different channels:

$$
\begin{equation*}
H=H_{B}+H_{b}+T_{\beta}+V_{\beta} \tag{4}
\end{equation*}
$$

$T$ and $V$ represents the relative motion and potential between two nuclei.
In general, the potential is a sum over all nucleons:

$$
\begin{equation*}
V_{\alpha}=\sum_{i \in a, j \in A} V_{i j} \quad \text { or } \quad V_{\beta}=\sum_{i \in b, j \in B} V_{i j} \tag{5}
\end{equation*}
$$

For convenience, we denote the total intrinsic Hamiltonian of $\alpha$ channel as:

$$
\begin{equation*}
H_{\alpha}=H_{A}+H_{a} \tag{6}
\end{equation*}
$$

So the total intrinsic state can be described by:

$$
\begin{equation*}
H_{\alpha} \Phi_{\alpha}=\epsilon_{\alpha} \Phi_{\alpha} \tag{7}
\end{equation*}
$$

## 2 Integral equation

The Schrodinger equation for the whole system is:

$$
\begin{equation*}
(E-H) \Psi=0 \tag{8}
\end{equation*}
$$

In a particular channel, this equation reads:

$$
\begin{equation*}
\left(E-H_{\alpha}-T_{\alpha}\right) \Psi_{\alpha}=V_{\alpha} \Psi_{\alpha} \tag{9}
\end{equation*}
$$

We then consider the free solution:

$$
\begin{equation*}
\left(E-H_{\alpha}-T_{\alpha}\right) \phi_{\alpha}=0 \tag{10}
\end{equation*}
$$

this equation can be solved through the separation of variables, so the solution takes the form:

$$
\begin{equation*}
\phi_{\alpha}=\exp \left(i \boldsymbol{k}_{\boldsymbol{\alpha}} \cdot \boldsymbol{r}_{\boldsymbol{\alpha}}\right) \times \Phi_{\alpha} \tag{11}
\end{equation*}
$$

where

$$
\begin{equation*}
\frac{\hbar^{2} k_{\alpha}^{2}}{2 \mu}=E_{\alpha}=E-\epsilon_{\alpha} \tag{12}
\end{equation*}
$$

is the kinetic energy of the relative motion. We introduce two different inner product notations:

$$
\begin{gather*}
\left(\Phi_{\alpha}, \Psi_{\alpha}\right)=\int \Phi_{\alpha}^{*} \Psi_{\alpha} \mathrm{d}(\text { internal coordinates })  \tag{13}\\
\left\langle\phi_{\alpha} \mid \Psi_{\alpha}\right\rangle=\int \phi_{\alpha}^{*} \Psi_{\alpha} \mathrm{d}(\text { all coordinates }) \tag{14}
\end{gather*}
$$

Since $\Psi_{\alpha}$ represents the state of the whole system, the relative motion state can be expressed by:

$$
\begin{equation*}
\psi_{\alpha}=\left(\Phi_{\alpha}, \Psi_{\alpha}\right) \tag{15}
\end{equation*}
$$

this is because that all the internal coordinates are integrated out. Project the total state onto intrinsic state of $\alpha$ channel, the Schrodinger equation then reads:

$$
\begin{equation*}
\left(E_{\alpha}-T_{\alpha}\right)\left(\Phi_{\alpha}, \Psi_{\alpha}\right)=\left(\Phi_{\alpha}, V_{\alpha} \Psi_{\alpha}\right) \tag{16}
\end{equation*}
$$

The equation of other channels can be obtained straightforward:

$$
\begin{equation*}
\left(E_{\beta}-T_{\beta}\right)\left(\Phi_{\beta}, \Psi_{\alpha}\right)=\left(\Phi_{\beta}, V_{\beta} \Psi_{\alpha}\right) \tag{17}
\end{equation*}
$$

The reason why $\Psi_{\alpha}$ is still denoted with $\alpha$ is that the incident channel is represented by $\alpha$.

### 2.1 T-matrix

Our goal is to solve the Schrodinger equation of the relative motion:

$$
\begin{equation*}
\left(E-H_{\alpha}-T_{\alpha}\right) \psi_{\alpha}=\left(\Phi_{\alpha}, V_{\alpha} \Psi_{\alpha}\right) \tag{18}
\end{equation*}
$$

So we define the Green's function:

$$
\begin{equation*}
\left(E_{\alpha}-T_{\alpha}\right) G_{\alpha}^{0}\left(\boldsymbol{r}_{\boldsymbol{\alpha}}-\boldsymbol{r}_{\boldsymbol{\alpha}}^{\prime}\right)=\delta\left(\boldsymbol{r}_{\boldsymbol{\alpha}}-\boldsymbol{r}_{\boldsymbol{\alpha}^{\prime}}\right) \tag{19}
\end{equation*}
$$

the asymptotic solution of Green's function is:

$$
\begin{equation*}
G_{\alpha}^{0(+)}\left(\boldsymbol{r}, \boldsymbol{r}^{\prime}\right)=-\frac{2 \mu}{4 \pi \hbar^{2}} \frac{e^{i k_{\alpha} r}}{r} \exp \left(-i \boldsymbol{k}_{\boldsymbol{\alpha}}^{\prime} \cdot \boldsymbol{r}^{\prime}\right) \tag{20}
\end{equation*}
$$

the plus sign represents the little positive imaginary part of in the denominator, which stands for the outgoing spherical solution.

The scattering amplitude can be expressed as:

$$
\begin{equation*}
f_{\alpha \alpha}(\theta)=-\frac{2 \mu_{\alpha}}{4 \pi \hbar^{2}} \int \mathrm{~d}^{3} r^{\prime} \exp \left(-i \boldsymbol{k}_{\alpha}^{\prime} \cdot \boldsymbol{r}^{\prime}\right)\left(\Phi_{\alpha}, V_{\alpha} \Psi_{\alpha}^{(+)}\right) \tag{21}
\end{equation*}
$$

similarly, the scattering amplitude for other channels is:

$$
\begin{equation*}
f_{\beta \alpha}(\theta)=-\frac{2 \mu_{\beta}}{4 \pi \hbar^{2}} \int \mathrm{~d}^{3} r^{\prime} \exp \left(-i \boldsymbol{k}_{\boldsymbol{\beta}}^{\prime} \cdot \boldsymbol{r}^{\prime}\right)\left(\Phi_{\beta}, V_{\beta} \Psi_{\alpha}^{(+)}\right) \tag{22}
\end{equation*}
$$

Finally we can express the scattering amplitude in a more general and compact way, by defining the T-matrix:

$$
\begin{equation*}
\mathscr{T}_{\beta \alpha}=\left\langle\phi_{\beta}\right| V_{\beta}\left|\Psi_{\alpha}^{(+)}\right\rangle \tag{23}
\end{equation*}
$$

the bracket notation is used to denote integration over all coordinates. $\beta$ stands for any specific channel. Remember that $\phi_{\text {beta }}$ is the free solution and $\Psi_{\alpha}^{(+)}$is the state vector of the whole system.

### 2.2 Asymptotic form of the complete wave function

Considering the elastic, inelastic and other break up channels, the complete wave function is written as:

$$
\begin{align*}
\Psi_{\alpha}^{(+)}= & \Phi_{\alpha}\left[\exp \left(i \boldsymbol{k}_{\boldsymbol{\alpha}} \cdot \boldsymbol{r}\right)+f_{\alpha \alpha}(\theta) \frac{e^{i k_{\alpha} r_{\alpha}}}{r_{\alpha}}\right] \\
& +\sum_{\alpha^{\prime} \neq \alpha} \Phi_{\alpha} f_{\alpha^{\prime} \alpha}(\theta) \frac{e^{i k_{\alpha^{\prime}} r_{\alpha}}}{r_{\alpha}}+\sum_{\beta} \Phi_{\beta} f_{\beta \alpha}(\theta) \frac{e^{i k_{\beta} r_{\beta}}}{r_{\beta}} \tag{24}
\end{align*}
$$

## 3 Formalism

### 3.1 LS equation

The Schrodinger equation reads:

$$
\begin{equation*}
\left(E-H_{\alpha}-T_{\alpha}\right)\left|\Psi_{\alpha}\right\rangle=V_{\alpha}\left|\Psi_{\alpha}\right\rangle \tag{25}
\end{equation*}
$$

including the free solution for boundary condition, we obtain the LS equation:

$$
\begin{equation*}
\left|\Psi_{\alpha}^{(+)}\right\rangle=\left|\phi_{\alpha}\right\rangle+\frac{1}{E-H_{\alpha}-T_{\alpha}+i \epsilon} V_{\alpha}\left|\Psi_{\alpha}^{(+)}\right\rangle \tag{26}
\end{equation*}
$$

similarly, we define:

$$
\begin{equation*}
\left|\Psi_{\alpha}^{(-)}\right\rangle=\left|\phi_{\alpha}\right\rangle+\frac{1}{E-H_{\alpha}-T_{\alpha}-i \epsilon} V_{\alpha}\left|\Psi_{\alpha}^{(-)}\right\rangle \tag{27}
\end{equation*}
$$

with the complete Hamiltonian:

$$
\begin{equation*}
\left|\Psi_{\alpha}^{(+)}\right\rangle=\left|\phi_{\alpha}\right\rangle+\frac{1}{E-H+i \epsilon} V_{\alpha}\left|\phi_{\alpha}\right\rangle \tag{28}
\end{equation*}
$$

this can be obtained with the identity:

$$
\begin{equation*}
\frac{1}{A}=\frac{1}{B}+\frac{1}{B}(B-A) \frac{1}{A} \tag{29}
\end{equation*}
$$

the Moller operator is defined as:

$$
\begin{equation*}
\left|\Psi_{\alpha}^{(+)}\right\rangle=\left(1+\frac{1}{E-H+i \epsilon} V_{\alpha}\right)\left|\phi_{\alpha}\right\rangle=\Omega_{\alpha}^{(+)}\left|\phi_{\alpha}\right\rangle \tag{30}
\end{equation*}
$$

### 3.2 T-matrix

The T-matrix is defined as:

$$
\begin{equation*}
\mathscr{T}_{\beta \alpha}=\left\langle\phi_{\beta}\right| V_{\beta}\left|\Psi_{\alpha}^{(+)}\right\rangle \tag{31}
\end{equation*}
$$

express the state of the whole system with Green's operator:

$$
\begin{equation*}
\mathscr{T}_{\beta \alpha}=\left\langle\phi_{\beta}\right| V_{\beta}\left(1+\frac{1}{E-H+i \epsilon} V_{\alpha}\right)\left|\phi_{\alpha}\right\rangle \tag{32}
\end{equation*}
$$

here we are concerned about the expectation of $V_{\beta}$ or $V_{\alpha}$ :

$$
\begin{align*}
\left\langle\phi_{\beta}\right| V_{\beta}-V_{\alpha}\left|\phi_{\alpha}\right\rangle & =\left\langle\phi_{\beta}\right|\left(H-H_{\beta}-T_{\beta}\right)-\left(H-H_{\alpha}-T_{\alpha}\right)\left|\phi_{\alpha}\right\rangle \\
& =\left\langle\phi_{\beta}\right|\left(H_{\alpha}+T_{\alpha}\right)-\left(H_{\beta}+T_{\beta}\right)\left|\phi_{\alpha}\right\rangle  \tag{33}\\
& =0
\end{align*}
$$

so we can replace $V_{\beta}$ with $V_{\alpha}$ :

$$
\begin{align*}
\mathscr{T}_{\beta \alpha} & =\left\langle\phi_{\beta}\right| V_{\alpha}+V_{\beta} \frac{1}{E-H+i \epsilon} V_{\alpha}\left|\phi_{\alpha}\right\rangle \\
& =\left\langle\phi_{\beta}\right|\left(1+V_{\beta} \frac{1}{E-H+i \epsilon}\right) V_{\alpha}\left|\phi_{\alpha}\right\rangle  \tag{34}\\
& =\left\langle\Psi_{\beta}^{(-)}\right| V_{\alpha}\left|\phi_{\alpha}\right\rangle
\end{align*}
$$

so we have two ways of writing the transition amplitude:

$$
\begin{equation*}
\mathscr{T}_{\beta \alpha}=\left\langle\phi_{\beta}\right| V_{\beta}\left|\Psi_{\alpha}^{(+)}\right\rangle=\left\langle\Psi_{\beta}^{(+)}\right| V_{\alpha}\left|\phi_{\alpha}\right\rangle \tag{35}
\end{equation*}
$$

these are sometimes referred as prior and post form of T-matrix.

